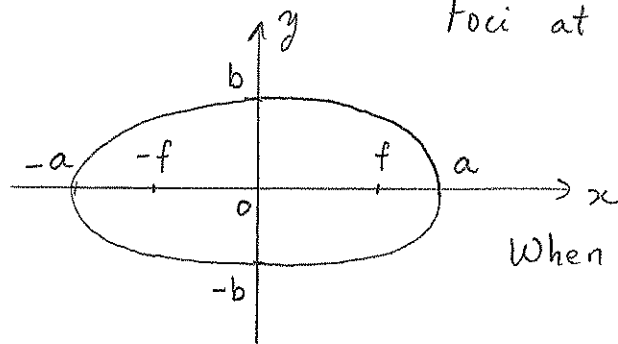


# CHAPTER 13

## Ellipses



Foci at  $\pm f$ .

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$

When  $a=b$  we get a circle.

Focus  $\rightarrow f = \sqrt{a^2 - b^2}$

Eccentricity  $\rightarrow e = \frac{\sqrt{a^2 - b^2}}{a} = \frac{f}{a}$

$a$  is semi-major axis length

$b$  is semi-minor axis length

Area  $\rightarrow \pi ab$

## Reasoning behind Newton's Law of Gravity

\* Centripetal acceleration of an object in circular motion

$$a_r = \frac{v^2}{r} \rightarrow \text{For any object on earth.}$$

Assuming same laws also apply to moon,

$$a_{\text{moon}} = g_{\text{at moon}} = \frac{v_m^2}{r_m}$$

Here  $v_m = 2\pi r_m / T_m \rightarrow 27.3 \text{ days}$

$$g_{\text{at moon}} = \frac{4\pi^2 r_m}{T_m^2}$$

\* From kinematics experiments with dropped objects on earth;

$$g_{\text{on earth}} = 9.80 \text{ m/s}^2.$$

\* Then, Newton looked at the ratio  $g_{\text{at moon}} / g_{\text{on earth}}$

$$\frac{g_{\text{at moon}}}{g_{\text{on earth}}} \approx \frac{1}{3600}$$

\* He also realized that  $\frac{r_m}{R_{\text{earth}}} \approx 60$  and  $60^2 = 3600$ .

$$\frac{g_{\text{at moon}}}{g_{\text{on earth}}} = \left( \frac{R_{\text{earth}}}{r_m} \right)^2$$

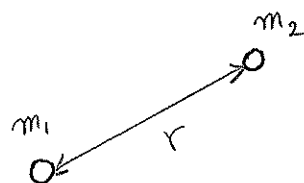
OR

$$\boxed{\frac{g_{\text{at moon}}}{g_{\text{on earth}}} = \left( \frac{1/r_m}{1/R_{\text{earth}}} \right)^2} \rightarrow \text{Inverse square Law!}$$

\* The key insight here was that same laws of physics govern the motion of planets, moon, etc. as those on earth. That is, laws of physics are universal.

Newton's Law of Gravity

$$F_{1\text{on}2} = F_{2\text{on}1} = \frac{G m_1 m_2}{r^2}$$



## Principle of Equivalence

\* Inertial mass of an object:

$$m_{\text{inertial}} = \frac{F}{a} \rightarrow \text{This has nothing to do w/ gravity.}$$

\* Gravitational mass of an object:

$$m_{\text{grav}} = \frac{r^2 F_{M \text{ on } m}}{GM} \quad \text{from } F_{M \text{ on } m} = \frac{GMm}{r^2}$$

\*  $m_{\text{inertial}} = m_{\text{grav}}$  is called principle of equivalence.

## Newton's Theory of Gravity

1) Newton's Law of Gravity:  $F_{m \text{ on } M} = \frac{GMm}{r^2}$

2) The principle of Equivalence.

3) Laws of motion are universal.

Little  $g$  and Big  $G$

$$\begin{aligned} & * \text{ On planet } X \rightarrow F_{g \text{ on } x} = m g_{\text{on } x} \\ & * \text{ Also } \rightarrow F_{M_x \text{ on } m} = \frac{GM_x m}{R^2} \end{aligned} \left. \vphantom{\begin{aligned} & * \text{ On planet } X \rightarrow F_{g \text{ on } x} = m g_{\text{on } x} \\ & * \text{ Also } \rightarrow F_{M_x \text{ on } m} = \frac{GM_x m}{R^2} \end{aligned}} \right\} g_{\text{on } x} = \frac{GM_x}{R^2}$$

\* Note that  $g$  decreases with distance as  $1/r^2$ .

Example

A planet has 4 times the mass of the earth, but the acceleration due to gravity on the planet's surface is the same as on the earth's surface. What is the planet's radius?

$$g_{\text{earth}} = \frac{GM_{\text{earth}}}{R_e^2} \quad \text{and} \quad g_x = \frac{GM_x}{R_x^2}; \quad M_x = 4M_{\text{earth}}$$

$$g_{\text{earth}} = g_x$$

$$\frac{GM_{\text{earth}}}{R_e^2} = \frac{4GM_{\text{earth}}}{R_x^2}$$

$$R_x^2 = 4R_e^2 \rightarrow$$

$$R_x = 2R_e$$

Gravitational Potential Energy.

$$\Delta U = -W_c(i \rightarrow f) = - \int_{x_i}^{x_f} F_x dx$$

\* Before, we used

$$F = -mg \quad \text{and} \quad \text{choose} \quad U = 0 \quad @ \quad y = 0.$$

to get  $U_g = mgy$ .  $\rightarrow$  This is true in the flat earth approximation when  $g$  is constant.

$$\Delta U = \underbrace{U_{\text{at } \infty}}_{\text{Choose to be 0.}} - U_{\text{at } r} = - \int_r^{\infty} (F_{10n2})_x dx$$

$$-U_{at r} = - \int_r^{\infty} - \frac{GMm}{x^2} dx = GMm \int_r^{\infty} \frac{dx}{x^2}$$

$$-U_{at r} = -GMm \left. \frac{1}{x} \right|_r^{\infty}$$

$$U_{at r} \equiv \boxed{U_g = -\frac{GMm}{r}}$$

We integrated  $F$  along the  $x$ -axis but since  $F$  is conservative, this holds for any path.

Example

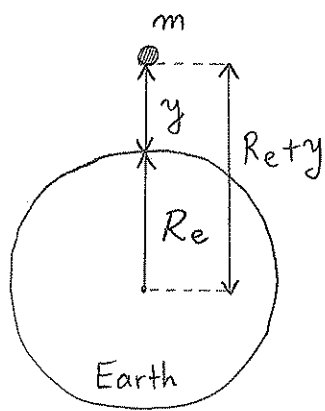
What is the "escape speed" for a 1000 kg rocket from the earth?

Energy Conservation:  $\underbrace{K_f + U_{g_f}}_{@ r \rightarrow \infty} = K_i + U_{g_i}$

$$0 + 0 = \frac{1}{2} m v_i^2 - \frac{GM_e m}{R_e}$$

$$v_i = \left( \frac{2GM_e}{R_e} \right)^{1/2} \approx 25,000 \text{ mph.}$$

## The Flat Earth Approximation



$$U_g = - \frac{GM_e m}{r} = - \frac{GM_e m}{R_e + y}$$

$$U_g = - \frac{GM_e m}{R_e (1 + y/R_e)}$$

Using the binomial approximation:  
 $(1+x)^n \approx 1 + nx$  if  $x \ll 1$ .

$$U_g \approx - GM_e m \frac{1}{R_e} (1 - y/R_e)$$

$$U_g \approx - \underbrace{\frac{GM_e m}{R_e}}_{\text{Gravitational Potential @ surface of earth.} \equiv U_0} + m \underbrace{\left[ \frac{GM_e}{R_e^2} \right]}_{= g_e} y \rightarrow \boxed{U_g = U_0 + mg_e y \text{ if } y \ll R_e}$$

Flat - Earth Approx.

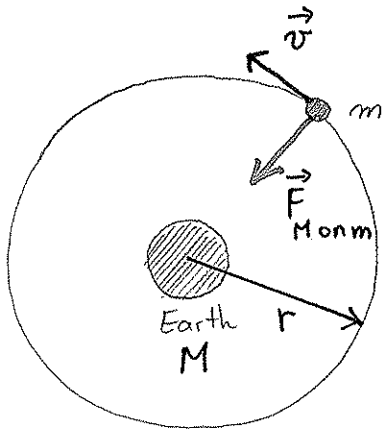
\* If we redefine our reference point for  $U_g$  to 0 rather than  $U_0$ ,

$$\boxed{U_g = mg_e y} \quad \text{Flat - Earth approximation.}$$

## 13.6 Satellite Orbits and Energies

Our analysis will be for circular orbits.

For earth,  $\frac{\text{semi-minor axis}}{\text{semi-major axis}} \approx 0.99986 \rightarrow$  very close to being circular.



$$F_{Monm} = \frac{GMm}{r^2} = \frac{mv^2}{r} = ma_r$$

Then speed of a satellite is

$$\frac{GMm}{r^2} = \frac{mv^2}{r}$$

$$v = \sqrt{\frac{GM}{r}}$$

→ Independent of  $m$ !

\*  $m$  is in free fall towards the earth

Example

The space shuttle in a 300 km high orbit ( $\approx 180$  mi) wants to capture a smaller satellite for repair. What are the speeds of the shuttle and the satellite in the orbit?

$$v_{\text{shuttle}} = v_{\text{satellite}} = v = \sqrt{\frac{(6.67 \times 10^{-11} \text{ Nm}^2/\text{kg}^2)(5.98 \times 10^{24} \text{ kg})}{\underbrace{3.0 \times 10^5 \text{ m}}_h + \underbrace{6.37 \times 10^6 \text{ m}}_{R_e}}}$$

$$v = 7730 \text{ m/s} \approx 17,000 \text{ mph.}^h$$

Kepler's Third Law

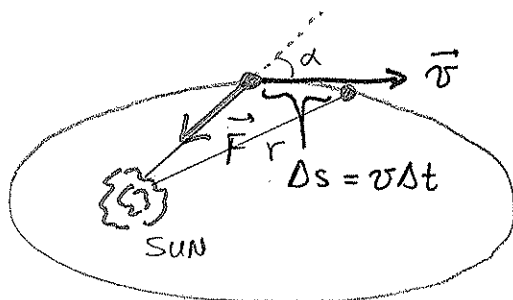
Period of the orbit is  $T \rightarrow v = \frac{2\pi r}{T}$

For a circular orbit  $\rightarrow v = \sqrt{\frac{GM}{r}} = \frac{2\pi r}{T}$

$$T^2 = \left( \frac{4\pi^2}{GM} \right) r^3$$

\* Square of the period is proportional to the cube of the radius. This is Kepler's 3rd Law.

## Kepler's 2nd Law

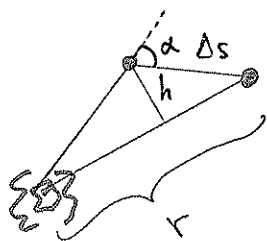


In small time interval  $\Delta t$ , the area swept is,

$$\Delta A = \frac{1}{2} \times \text{base} \times \text{height}$$

$$\Delta A = \frac{1}{2} r \Delta s \sin \alpha$$

$$= v \Delta t$$



$$h = \Delta s \sin \alpha$$

$$\frac{\Delta A}{\Delta t} = \frac{1}{2} r v \sin \alpha \cdot \frac{m}{m}$$

$$\frac{\Delta A}{\Delta t} = \frac{m v r \sin \alpha}{2m} = \frac{L}{2m}$$

The angular momentum  $L$  is conserved since  $\tau_{\text{net}} = 0$ ,

$$\boxed{\frac{\Delta A}{\Delta t} = \text{constant}} \quad \text{Kepler's 2nd Law}$$

## Kinetic Energy vs. Potential Energy in Orbit

$$\left. \begin{array}{l} \text{A satellite's KE} \rightarrow K = \frac{1}{2} m v^2 \\ \text{For a circular orbit} \rightarrow v = \sqrt{\frac{GM}{r}} \end{array} \right\} K = \frac{1}{2} m \frac{GM}{r}$$

$$\boxed{K = \frac{1}{2} \underbrace{\frac{GMm}{r}}_{-U_g} = -\frac{1}{2} U_g}$$

For a circular orbit.



$$E_{\text{mech}} = K + U_g$$

$$E_{\text{mech}} = -\frac{1}{2} U_g + U_g = \frac{1}{2} U_g$$

\* Note that  $U_g < 0$  hence  $E_{\text{mech}} < 0$ . This is characteristic of a "bound" system.

\*  $E_{\text{mech}}$  for an unbounded system must be positive since  $U \rightarrow 0$  as  $r \rightarrow \infty$  and  $K \geq 0$  always.