Assessing tunneling-multiphoton dichotomy in photoionization: Keldysh parameter γ versus scaled frequency Ω

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Motivation and previous work

• Keldysh parameter $\gamma = \omega \sqrt{2I_p}/F$ is commonly used for referring to the dominant ionization regime in atoms: $\gamma < 1$ for tunneling and $\gamma > 1$ for multiphoton ionization.

[L. V. Keldysh, Zh. Eksp. Theor. Fiz. 47, 1945 (1964)]

- There are several shortcomings of this approach when laser frequencies are considered outside the usual strong field process driven by 800 nm laser pulses, *e.g.* HHG and ATI. [H. R. Reiss, Prog. Quant. Electr. 16, 1 (1992).]
- For a fixed laser frequency, laser intensity can be varied so that **any** value for γ can be attained.
- We map out a landscape in (γ, Ω) space, which shows regions with features commonly attributed to tunneling or multiphoton processes.
- Expressing γ in terms of the scaled parameters $\Omega = \omega n^3$ and F/n^4 gives $\gamma = \Omega/F$, which takes dynamics inside the Coulomb potential into account in addition to the laser induced dynamics.
- We perform two sets of calculations for ionization rates and probabilities of a H atom out of 1s, 2s, 8s and 16s states for a large set of (γ, Ω) pairs.

Numerical Simulations

We solve the time-dependent Schrödinger equation with a linearly polarized laser field on a (ℓ, r) -grid using an implicit split split operator technique.

Ionization Rate Calculations

We express the total wave function as a superposition $\psi(r, l, t) = \psi_0(r, l) + \psi_1(r, l, t)$, and solve

$$\left[i\frac{\partial}{\partial t} - H(r,l,t)\right]\psi_1(r,l,t) = H_{\mathsf{L}}\psi_0$$

Given that $\psi(t) = \psi_0 + \psi_1(t)$, and that ψ_0 is an eigenstate, this equation is exact and allows for atomic processes of all orders.

Ionization Probability Calculations

We solve the full time-dependent Schrödinger equation starting from the initial states 1s, 4s, and 8s for a laser pulse with width of 160 classical periods.

 $\left[i\frac{\partial}{\partial t} - H(r,l,t)\right]\psi(r,l,t) = 0$

Ionization Rates in (γ, Ω) -space

Ionization rates Γ from initial states 1s, 4s, 8s, and 16s. The Ω and Γ axes are in \log_2 and \log_{10} scale. The ridge structures correspond to the N-photon ionization thresholds at



Constant F and Ω slices

Slices along constant F and Ω lines in the (γ, Ω) landscape yield the rates as function of Ω (spectroscopic viewpoint) and F (strong field viewpoint).



(Left column) Ionization probabilities out of 1s. 4s, points. Momenta corresponding to lowest order (diand 8s states for large $\gamma \gg 1$ and Ω up to 4 using a rect) multiphoton ionization, the 2U_p and 10U_p limits laser pulse with a fixed width of 160 τ_{Rvd} . $\gamma \gg 1$ im- are indicated by dashed semicircles. The plots are plies tunneling is predominant in all three maps. The in \log_2 (for 1s) and \log_4 scales (for 4s and 8s). Multimultiphoton peaks and suppression of ionization are photon rings and tunneling (small p) regions are visiclearly visible in the $\Omega < n/2$ and $\Omega > n/2$ regions. (Right columns) Momentum maps at select few rescattering rings.







The common feature in all of the momentum maps is that they are dominated by the low energy electrons a distinct sign of tunneling. Despite the high values of the Keldysh parameter, tunneling and multiphoton processes coexist in this regime.



Ionization Probability and Momentum Distributions

ble along with some well known structures, like back

Momentum distributions for $\Omega = 8$. Although Keldysh theory predicts tunneling to dominate the ionization dynamics in these cases (γ <1), momentum distribution is concentrated on the circular ridge with the radius suggesting single-photon ionization. The ridge is sharp for the higher γ (lower intensity), and dispersed for the lower γ (higher intensity). No electrons inside the 10 U_p limit for $\gamma = 0.8$ meaning no ionized electrons are due to tunneling. Also no electrons for $E < 2U_p$ for $\gamma = 0.4$.

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